

# Statistical Orbit Determination



Lecture 5 – Perturbed Motion  
Presenter: Christopher R. Simpson

# Recap

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- Lecture 4 – Notes posted [here](#)
  - Classical two-body problem
- Lecture 6 – Coordinate systems and time
- Questions
  - Post them to YouTube page



# Agenda

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- Perturbed Motion
  - Classical example: Lunar problem
  - Variation of parameters
  - Gravitational perturbations
    - Oblateness
    - Third-body Effects
  - Nongravitational perturbations
- Assigned Problems

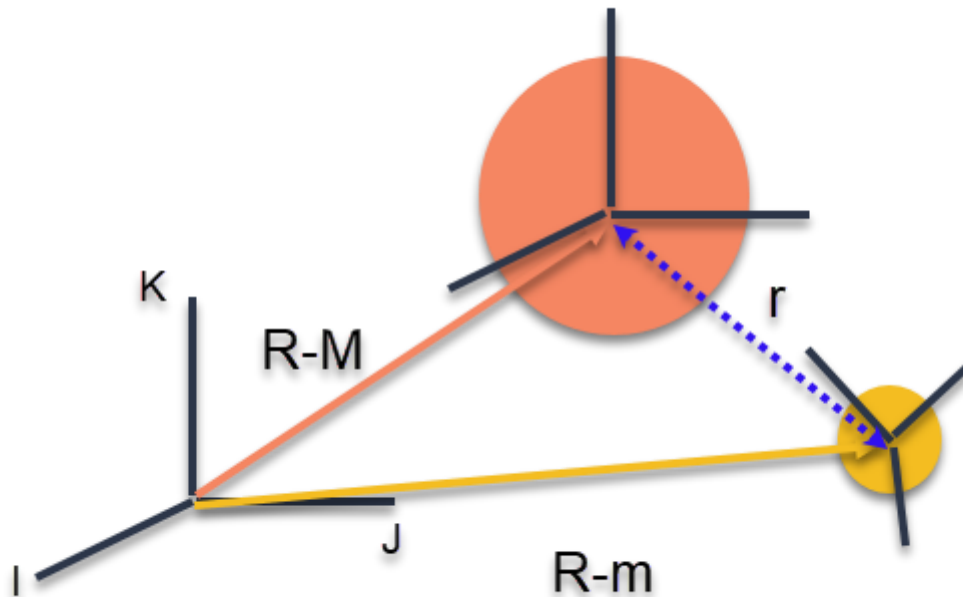


# Perturbed Motion – Recap of two body (1/2)

- Relative motion of  $m_2$  w/r to  $m_1$

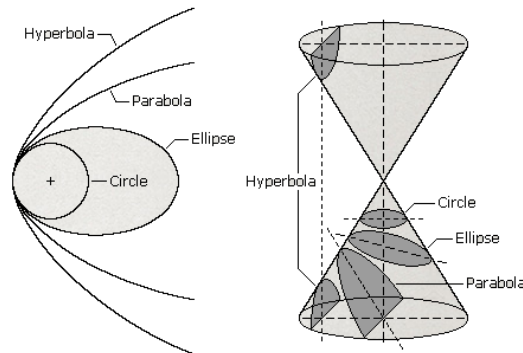
$$\ddot{\vec{r}} + \frac{\mu}{r^3} \vec{r} = 0$$

- Assumptions
  - Two point masses/bodies are spherically symmetric
  - Gravitational force propagates instantaneously (No relativistic effects)



# Perturbed Motion – Recap of two body (2/2)

- Previous description of motion is idealized
  - Equations of motion can be solved analytically
  - Motion simplified to a geometric shape (circle, ellipse, parabola, hyperbola)
- Newton
  - Told Halley, that the motion of the Moon [in the three body system] “made his head ache and kept him awake so often that he would think of it no more,” (Moulton, p. 363, 1914).
- No general closed-form solution for three-body problem
  - Approximate analytical solutions use two-bodies as a reference
  - Approximate or general perturbations solution adds contributing perturbing forces
- Numerical solutions
  - Perturbed motion represented by a set of ODEs with specified initial conditions (special perturbations)



# Perturbed Motion – Lunar problem (1/2)

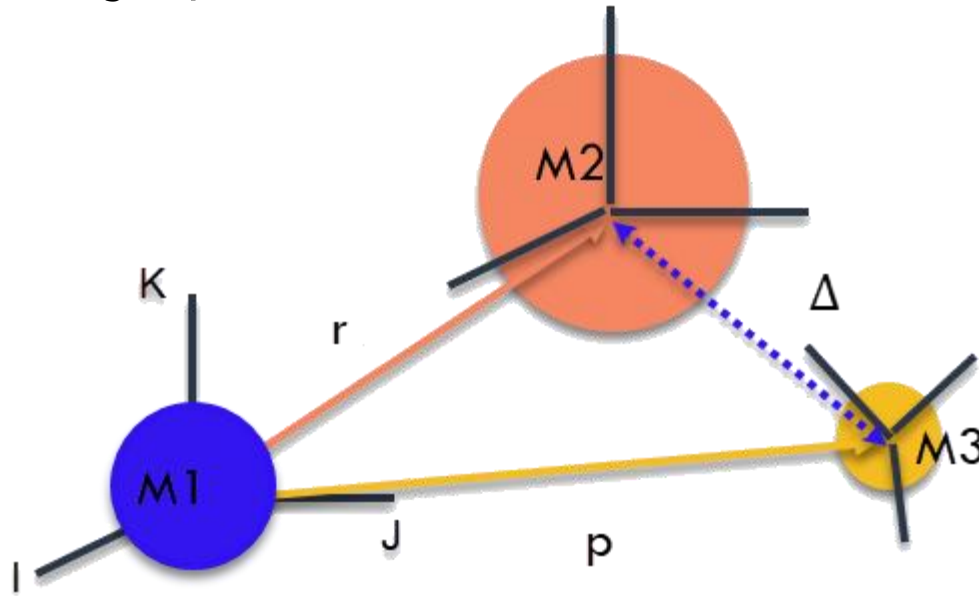
- Derive equations of motion for three body problem

- Solve EoM using numerical integration

$$\ddot{\vec{r}} = -\frac{\mu \vec{r}}{r^3} + GM_3 \left( \frac{\bar{\Delta}}{\Delta^3} - \frac{\vec{r}_p}{r_p^3} \right)$$
$$\ddot{\vec{r}}_p = -\frac{\mu' \vec{r}}{r^3} + GM_2 \left( \frac{\bar{\Delta}}{\Delta^3} + \frac{\vec{r}}{r^3} \right)$$

- where  $\mu = G(M_1 + M_2)$  and  $\mu' = G(M_1 + M_3)$

- Let  $M_1$ ,  $M_2$ , and  $M_3$  represent Earth, the Moon, and the Sun, respectively



# Perturbed Motion – Lunar problem (2/2)

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- Restricted three-body problem (Szebehely, 1967)
  - Sun's mass is more than 300,000 times greater than Earth
  - Moon's mass is 81 times smaller than Earth
$$\mu = GM_1$$
$$\mu' = GM_3$$
- Can integrate w/o approximation given specified initial conditions
  - (Shampine and Gordon, 1975), since Sun is dominant perturbation of the lunar motion
  - Rotate about  $X$  in ECI frame so  $Z$  axis is perpendicular to ecliptic
  - Can show osculating elements (orbital elements) are not constant due to perturbations
- Osculating element variation
  - Ascending node linear variation with time (secular variation/periodic variations)
  - Inclination has no apparent secular but experiences periodic variation
  - Secular node rate is negative (regression of the node,  $\sim 19.4^\circ$  per year)



# Perturbed Motion – Variation of parameters

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- Temporal variations of Moon's osculating elements

$$\ddot{\vec{r}} = -\frac{\mu}{r^3} \vec{r} + \bar{\vec{f}}$$

- Develop solution to ODEs by using variation of parameters
  - $\bar{\vec{f}}$  is perturbing force
  - See Appendix D for differential equations describing change of osculating elements
- In some cases  $\bar{\vec{f}}$  is derivable from potential or disturbing function
- $\bar{\vec{f}}$  can be categorized as gravitational or nongravitational





# Gravitational – Mass Distribution (1/6)

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- Mass distribution

- Two point masses gravitational potential

$$U = \frac{GM_1M_2}{r}$$

$$\bar{F} = \nabla U = \frac{GM_1M_2}{r^3} \bar{r}$$

- Can model mass distribution as collection of point masses

- Potential experienced by point mass,  $m'$ , is

$$U = m' \int \int \int G\gamma \, dx \, dy \, dz / \rho$$

where  $\gamma$  is the mass density associated with  $dm$ ,  $dx \, dy \, dz$  are differential volume, and  $\rho$  is distance between differential mass and external mass  $m'$

- Spherical harmonics splits Earth into regions

- Allows us to assign mass coefficients/properties to each region



# Gravitational – Mass Distribution (2/6)

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- Taking the external mass to be unity,  $m' = 1$ 
  - $U = \int_M \frac{Gdm}{\rho}$ , where we are integrating over entire mass
  - Position vector of  $m'$  is  $\bar{r}$
  - Where  $(x, y, z)$  is considered to be body-fixed
  - $(x, y)$  equatorial plane and  $x$  to Greenwich meridian
- Expand using infinite series
  - Expand  $U = \int_M \frac{Gdm}{\rho}$  using infinite series and Legendre polynomials
$$U = \frac{G}{r} \int_M \sum_{l=0}^{\infty} \left(\frac{R}{r}\right)^l P_l(\cos S) dm$$
  - Where  $R$  is distance between origin and  $dm$  and  $P_l$  is Legendre polynomial of degree  $l$  with an argument equal to the cosine of the angle between the two vectors  $R$  and  $r$



# Gravitational – Mass Distribution (3/6)

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- Expand the Legendre polynomial into spherical harmonics
  - Terms dependent on mass distribution are collected into coefficients

$$U = \frac{\mu}{r} + U'$$

$$U' = -\frac{\mu^*}{r} \sum_{l=1}^{\infty} \left(\frac{a_e}{r}\right)^l P_l(\sin \phi) J_l + \frac{\mu^*}{r} \sum_{l=1}^{\infty} \sum_{m=1}^l \left(\frac{a_e}{r}\right)^l P_{lm}(\sin \phi) [C_{lm} \cos m\lambda + S_{lm} \sin m\lambda]$$

- Coordinates of  $m'$  are now expressed in spherical coordinates  $(r, \phi, \lambda)$ 
  - $\phi$  is geocentric latitude
  - $\lambda$  is longitude angle
- Scale factors to nondimensionalize  $C_{lm}$  and  $S_{lm}$ 
  - Reference mass,  $\mu^* = GM^*$
  - Reference distance,  $a_e$



# Gravitational – Mass Distribution (4/6)

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- Zonal harmonics ( $m = 0$ ),  $J_l$ 
  - No dependence on longitude
  - Circle of latitude alternately positive and negative
- Sectorial harmonics ( $n = m$ )
  - Sectors alternately positive and negative along lines of longitude
- Tesseral harmonics ( $n \neq m$ )
  - Checkerboard array of domains, “square” harmonics



Zonal Harmonics



Sectorial Harmonics



Tesseral Harmonics



# Gravitational – Mass Distribution (5/6)

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- Normalized expressions normally used
  - Legendre functions normally high numerical values compared to mass coefficients
- Degree 1 terms,  $l$ , proportional to distance between cm and  $O$
- Degree 2 terms proportional to moments and products of inertia

$$\begin{aligned}\bar{F}^* &= m' \nabla U \\ \ddot{\vec{r}} &= \left(1 + \frac{m'}{M}\right) \nabla U\end{aligned}$$

- If  $m'/M$  is very small

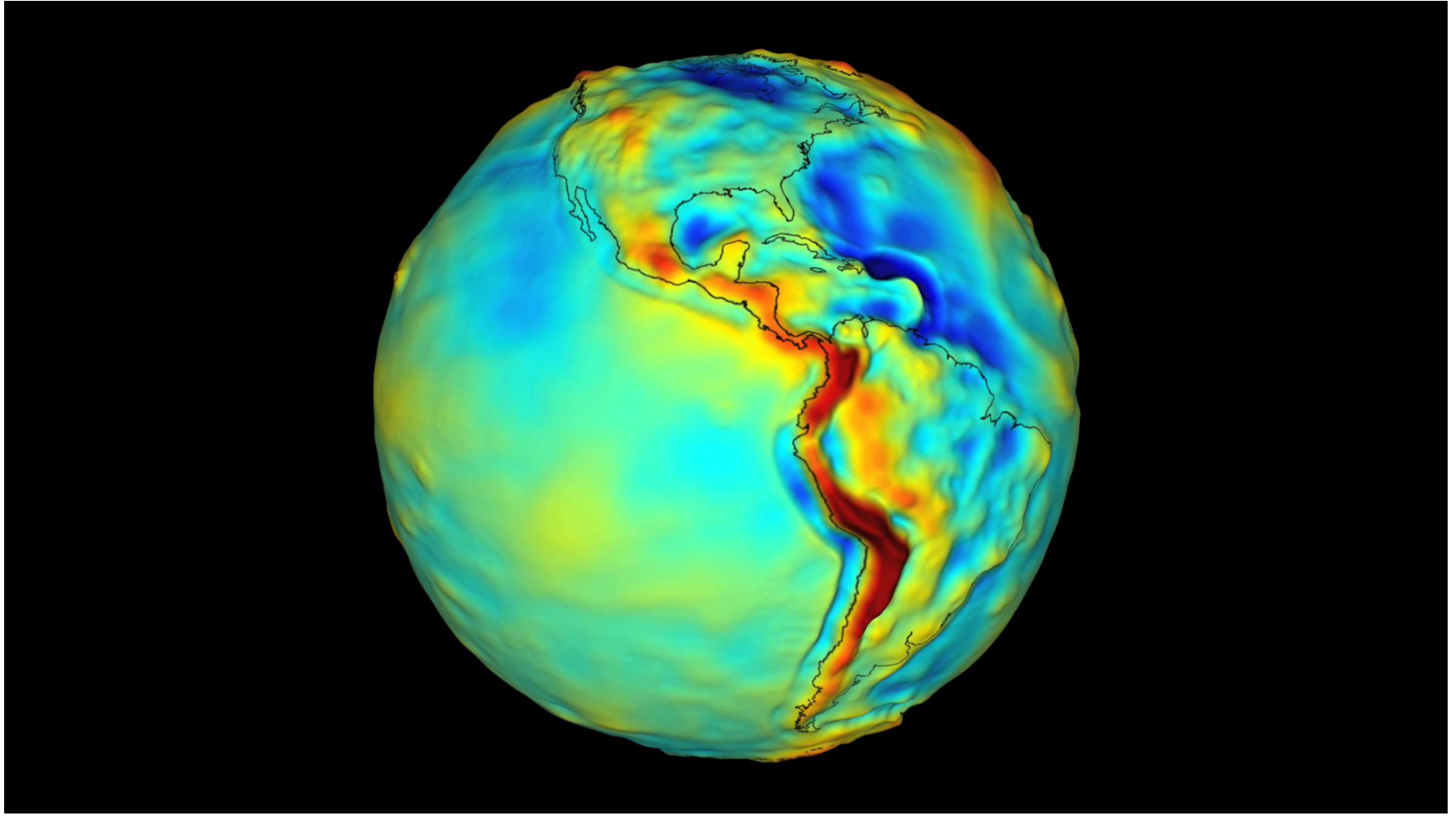
$$\ddot{\vec{r}} = \nabla U = -\frac{\mu \bar{r}}{r^3} + \bar{f}_{NS}$$

- Be careful about whether system used is nonrotating or rotating



# Gravitational – Mass Distribution (6/6)

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# Gravitational – Oblateness (1/3)

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- More than 95% of gravitational force (that's not  $\mu/r^2$ ) is  $J_2$

- Potential for the ellipsoid of revolution

$$U' = -\frac{\mu}{r} \left(\frac{a_e}{r}\right)^2 J_2 P_2(\sin \phi)$$

- Can relate to orbit elements

$$\sin \phi = \sin i \sin(\omega + v)$$

- Gravitational potential now can be expressed in terms of orbit elements

- Use eccentricity expansions

$$U' = -\frac{\mu}{a} \left(\frac{a_e}{a}\right)^2 J_2 \{3/4 \sin^2 i [1 - \cos(2\omega + 2M)] - 1/2\} + \text{higher order terms}$$

- Can divide contributions into secular and periodic

$$U' = U_s + U_p$$

- Secular

$$U_s = -\frac{GM}{a} \left(\frac{a_e}{a}\right)^2 J_2 \left(\frac{3}{4} \sin^2 i - \frac{1}{2}\right)$$

- Periodic

$$U_p = \frac{GM}{a} \left(\frac{a_e}{a}\right)^2 J_2 \left(\frac{3}{4} \sin^2 i \cos(2\omega + 2M)\right)$$



# Gravitational – Oblateness (2/3)

- Which orbit elements affected over time

- $a, e, i$  not affected by time
- $\dot{\Omega}_s$ , secular node rate is constant for given  $a, e, i$

$$\dot{\Omega}_s \cong -\frac{3}{2}J_2 n \left(\frac{a_e}{a}\right)^2 \cos i$$

- Application in solar-synchronous satellites

- Constantly aligned with Earth-Sun direction
- $\dot{\Omega}_s = 360^\circ/365.25 \text{ days} \cong 1^\circ/\text{day}$
- In other words, for  $e \approx 0$ ,  $a \approx 7000 \text{ km}$ , and  $i \approx 98$

- Secular rates of  $\Omega$ ,  $\omega$ , and  $M$  (Kaula, 1966)

$$\dot{\Omega}_s = -\frac{3}{2}J_2 \frac{n}{(1-e^2)^2} \left(\frac{a_e}{a}\right)^2 \cos i$$

$$\dot{\omega}_s = \frac{3}{4}J_2 \frac{n}{(1-e^2)^2} \left(\frac{a_e}{a}\right)^2 (5 \cos^2 i - 1)$$

$$\dot{M}_s = \bar{n} + \frac{3}{4}J_2 \frac{n}{(1-e^2)^{3/2}} \left(\frac{a_e}{a}\right)^2 (3 \cos^2 i - 1)$$





# Gravitational – Oblateness (3/3)

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- Previous equations (Kaula, 1966) use mean elements
  - $a, e, i$  have periodic variations averaged out
  - $\bar{n} = \sqrt{\mu/\bar{a}^3}$  where  $\bar{a}$  is the mean value
- Can use this for simple model, Secularly Precessing Ellipse
- From Kaula's linear derivation
  - Nodal period,  $P_n = \frac{2\pi}{\dot{\omega}_s + \dot{M}_s}$
  - Nodal day,  $D_n = \frac{2\pi}{\dot{\Omega}_s + \omega_e}$



# Gravitational – Third-body effects

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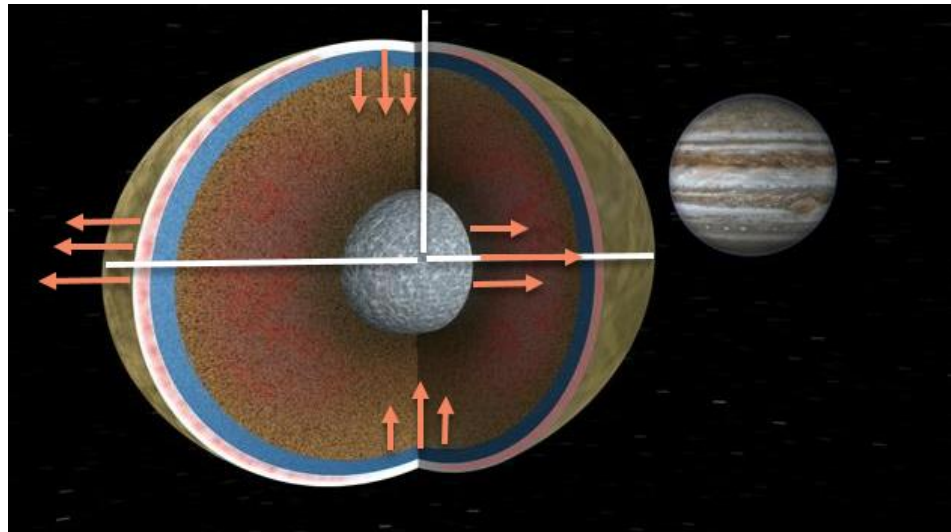
- Consider the two-body case, then add another body
  - Not considering the non-uniformity of the gravitational field just yet
  - Addition given by  $f_{3b} = \sum_{j=1}^{n_p} \mu_j \left( \frac{\bar{\Delta}_j}{\Delta_j^3} - \frac{\bar{r}_j}{r_j^3} \right)$
  - $\bar{\Delta}_j$  is the vector between the  $j$  body and the satellite
  - $\bar{r}_j$  is the vector between the  $j$  body and the Earth



# Gravitational – Tidal Effects (1/3)

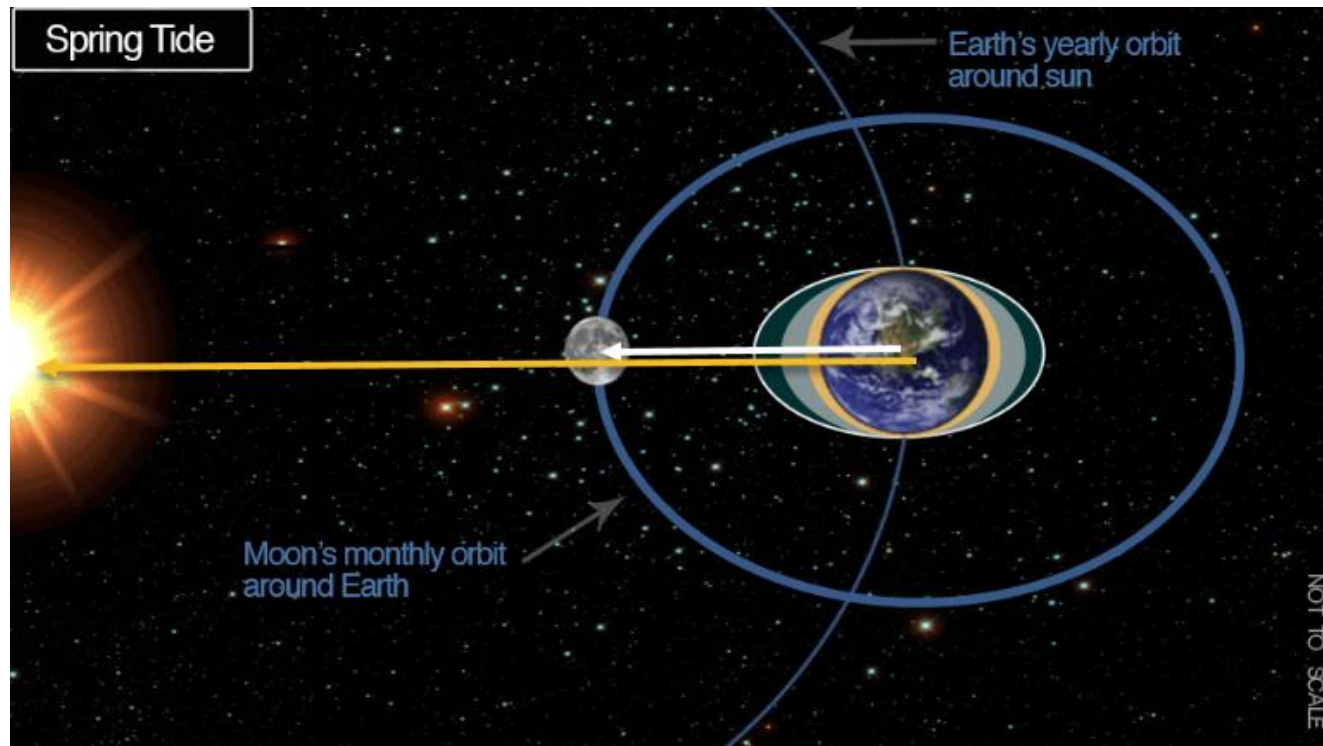
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- One side of celestial body experiences greater acceleration
  - Redistributes liquid and solid mass of body
  - Measured  $w/r$  to center of mass
- Rotation of body creates a phase advance
  - Accelerates moon
  - Tidal locking, moon's tidal bulge and face match Earth



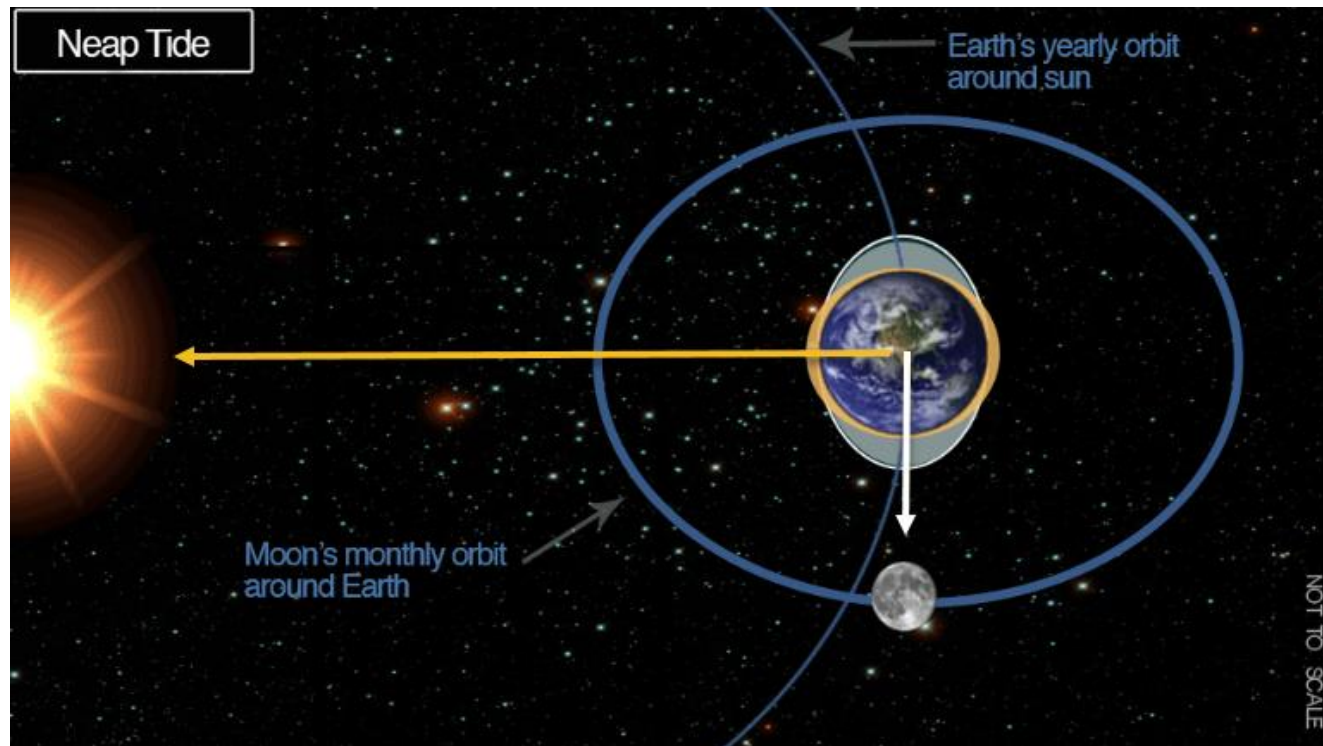
# Gravitational – Tidal Effects (2/3)

- Changes GS position and SC acceleration
- Sun and Moon tidal interaction
  - Spring Tides and Neap Tides



# Gravitational – Tidal Effects (2/3)

- Changes GS position and SC acceleration
- Sun and Moon tidal interaction
  - Spring Tides and Neap Tides



# Gravitational – Tidal Effects (3/3)

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- Tidal potential approximated by

$$U_t = -\frac{\mu_{\text{moon}}}{R} \sum_{n=2}^{\infty} k_n \left(\frac{r}{R}\right)^n P_n(\cos \psi)$$

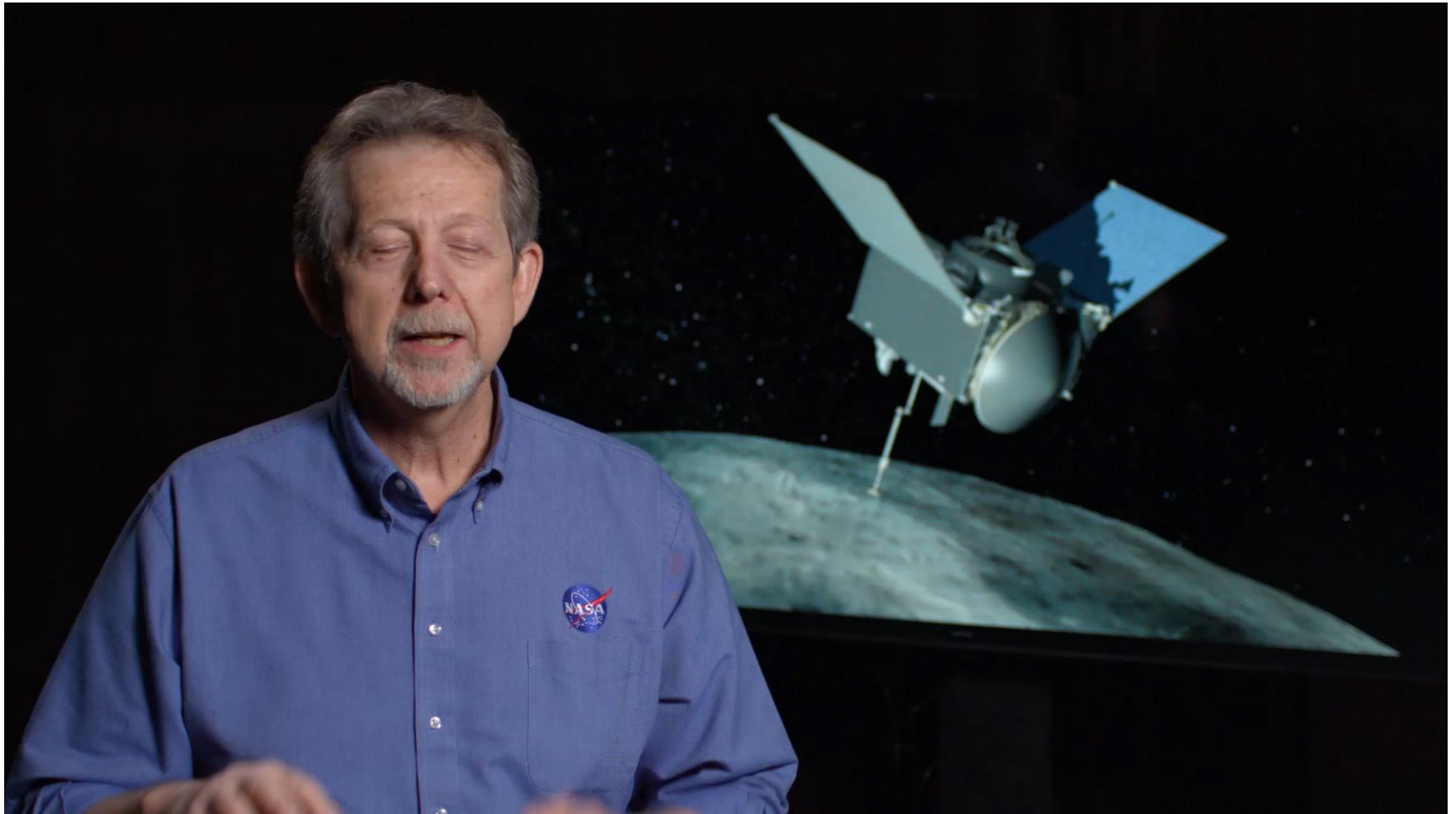
- Love numbers,  $k_n$ , describe how deformable celestial body is
  - $r, \psi$  describe position on affected body looking down on orbital plane
- Variations in topology change amplitudes of tides

Celestial Body	Ocean	Solid
Sun, ☉	0.82 ft	0.60 ft
Moon, ☾	1.8 ft	1.3 ft



# Applications – Lagrange Points (1/2)

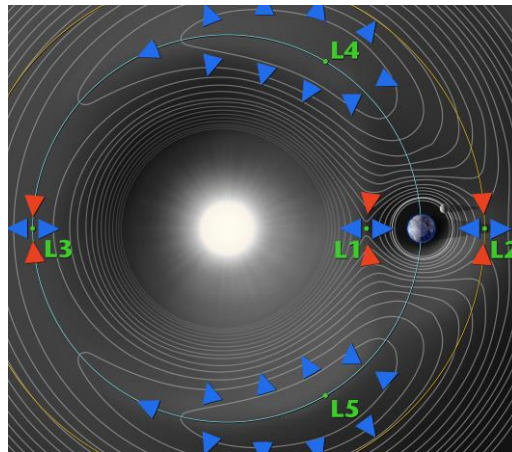
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# Applications – Lagrange Points (2/2)

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- Points where gravitational pull is balanced
  - Exists between all multibody systems
  - Sun-Earth Lagrange points most commonly used/considered
- Several current missions either use or will pass L points
  - OSIRIS-REx will swing by L4 looking for Earth Trojans
  - JWST will orbit L2 to examine early origins of universe
  - Solar and Heliospheric Observatory (SOHO) orbits L1





# Perturbed Motion – Nongravitational

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- Atmospheric Resistance (Drag)
  - $\bar{f}_D = -\frac{1}{2}\rho\left(\frac{C_DA}{m}\right)v\bar{v}$ , force of drag acting opposite to movement
  - Ballistic coefficient,  $\left(\frac{C_DA}{m}\right)$ 
    - At low altitudes ( $\sim 350$  km) the atmospheric density is  $10^{-11}$  of sea level
    - Mean free path increases to meters
  - Drag removes energy from orbit
    - Secular decay in semimajor axis and eccentricity
    - Decay in  $a$  will determine lifetime of satellite
- Some re-entry numbers
  - $\sim 170$  million pieces of space debris
  - Less than 0.2% tracked

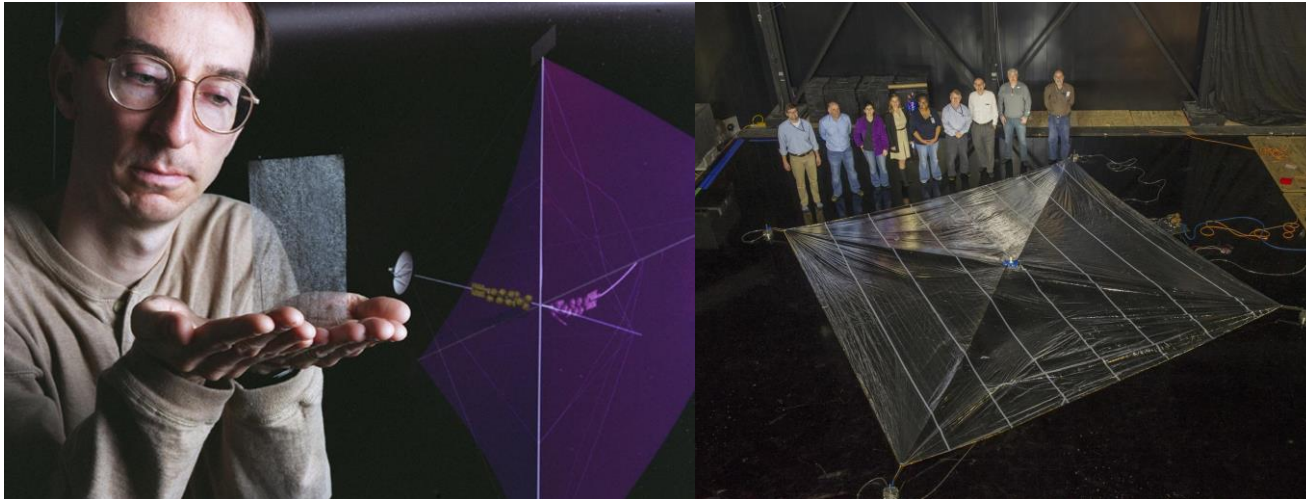


# Perturbed Motion – Nongravitational

- Solar radiation pressure
  - Transfer of momentum through photons

$$f_{SRP} = -P \frac{vA}{m} C_R \bar{u}$$

- $P$  is the momentum flux
- $A$  is cross-sectional area
- $C_R$ , reflectivity coefficient
- $v$ , eclipse factor for when satellite is in shadow





Practice problems: The Orbit Problem

# **PREDICT THE ORBIT**

# Assigned Problems - Overview

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- You are given three problems involving orbital motion. They have been picked to ensure you have a sufficient understanding of orbital mechanics before proceeding. The problems resemble numbers 4, 5, 6, 10, 11, and 12 from the textbook.
- These problems should be complete by Friday, February 8.



# Assigned Problems – Problem 1

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- Given the following position and velocity of a satellite
  - Expressed in a non-rotating geocentric coordinate system

	Position (m)	Velocity (m/s)
X	7088580.789	-10.20544809
Y	-64.326	-522.85385193
Z	920.514	7482.075141

- Determine the six orbital elements ( $a, e, i, \Omega, \omega, M_0$ )
- Assuming  $X_0$  is given and two-body motion, predict position and velocity at  $t = 3,000$  sec. Determine flight path angle at this time.
- Determine the latitude and longitude of the subsatellite point for  $t = 3,000$  sec if  $\alpha_G$  at  $t = 0$  is 0. Assume the  $Z$  axis of the nonrotating system is coincident with the  $z$  axis of the rotating system.



# Assigned Problems – Problem 2 (1/2)

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- Orbit of CRISTA-SPAS-2

- [Cryogenic Infrared Spectrometers and Telescopes for the Atmosphere](#)

The joint venture of DLR and NASA, the small free-flying satellite contains three telescopes, four spectrometers, and a GPS receiver on-board. It is deployed from the shuttle Discovery on STS-85 in August 1997. Using on-board navigation, the receiver measurements are processed in an Earth-centered, Earth-fixed coordinate system.

August 18, 1997		
GPS-T (hrs:min:sec)	00:00:0.000000	00:00:03.000000
x	3325396.441	3309747.175
y	5472597.483	5485240.159
z	-2057129.050	-2048664.333

August 19, 1997		
GPS-T (hrs:min:sec)	00:00:0.000000	00:00:03.000000
x	4389882.255	4402505.030
y	-4444406.953	-4428002.728
z	-2508462.520	-2515303.456



# Assigned Problems – Problem 2 (2/2)

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a) Demonstrate that the node location is not fixed in space and determine an approximate rate of node change (degrees/day) from these positions.

Compare the node rate with the value predicted by

$$\dot{\Omega} = -\frac{3}{2}J_2 \frac{n}{(1-e^2)^2} \left(\frac{a_e}{a}\right)^2 \cos i$$

b) Determine the inclination of CRISTA-SPAS-2 during the first 3-sec interval and the last 3-sec interval.

Comment: The position vectors determined by GPS in this case are influenced at the 100-meter level by Selective Availability, but the error does not significantly affect this problem.



# Assigned Problems – Problem 3 (1/2)

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- GLONASS

- [Russia's answer for American GPS](#)

Given a set of initial conditions for a high-altitude GLONASS satellite, numerically integrate the equations of motion for one day.

- a) Assuming the satellite is influenced by  $J_2$  only, derive the equations of motion in non-rotation coordinates. Assume the nonrotating  $Z$  axis coincides with the Earth-fixed  $z$  axis.
- b) During the integration, compute the Jacobi constant and the  $Z$  component of the angular momentum. Are these quantities constant?
- c) Plot the six orbital elements as a function of time.
- d) Identify features similar to and different from Fig. 2.3.5





## Assigned Problems – Problem 3 (2/2)

e) Compare the node rate predicted by

$$\dot{\Omega} = -\frac{3}{2}J_2 \frac{n}{(1-e^2)^2} \left(\frac{a_e}{a}\right)^2 \cos i$$

with a value estimated from (c).

f) Compare the amplitude of the semimajor axis periodic term with

$$a(t) = \bar{a} + 3\bar{n}\bar{a}J_2 \left(\frac{a_e}{\bar{a}}\right)^2 \sin^2 \frac{\bar{i}(\cos(2\omega + 2M))}{2\dot{\omega}_s + 2\dot{M}_s}$$

g) Plot the ground track. Does the ground track repeat after one day?

$a$	$e$	$i$	$\Omega$	$\omega$	$M_0$
25500.0 km	0.0015	63 deg	-60 deg	0 deg	0 deg

